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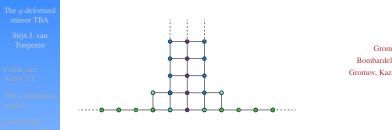
Stijn J. van Tongeren



Work done in collaboration with G. Arutyunov and M. de Leeuw, [1208.3478]



Introduction

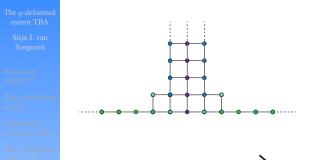


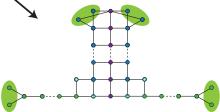
Arutyunov, Frolov '09 Gromov, Kazakov, Vieira '09 Bombardelli, Fioravanti, Tateo '09 Gromov, Kazakov, Kozak, Vieira '09

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• Interesting model



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► Integrability: trigonometric rather than rational



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 - ► Integrability: trigonometric rather than rational
 - ► TBA: interesting structure (XXZ)



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- ► (Thermodynamics of) the *q*-deformed Hubbard model



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Alcaraz and Bariev '99

• Conjuctured relation to Pohlmeyer reduced string theory



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 - q-deformed theory interpolates



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 - q = 1, g arbitrary: AdS₅ × S⁵ string theory
 - $q = e^{i\pi/k}, g \to \infty$: solitons of ssssG

Hoare and Tseytlin '11 Hoare, Hollowood and Miramontes '11



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 - Physical 'regularization' of the problem $(q = e^{i\pi/k})$



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- Complementary approach to $AdS_5\times S^5$ mirror TBA
 - Physical 'regularization' of the problem $(q = e^{i\pi/k})$
 - Wider perspective
- Possible (partial) applications to particular deformed backgrounds





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• TBA describes finite size string spectrum via a mirror model

Zamolodchikov '89 Ambjorn, Janik and Kristjansen '05 Arutyunov and Frolov '07



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• Today: the quantum deformation of this story (at roots of unity)



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 - ► Thermodynamic limit: different string hypothesis



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Conclusion

• The $\mathfrak{su}(2|2)$ superalgebra in Chevalley-Serre basis $(3 \times E, F, H)$ $[H_i, H_j] = 0, \ [H_i, E/F_j] = \pm A_{ij}E/F_j, \ [E_i, F_j] = \delta_{ij}D_iH_i,$

with

$$A = \begin{pmatrix} 2 & -1 & 0 \\ -1 & 0 & 1 \\ 0 & 1 & -2 \end{pmatrix}, D = \operatorname{diag}(1, -1, -1)$$

plus Serre relations



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plus (deformed) Serre relations



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• We take $q = e^{i\pi/k}$ with integer k > 2



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• q-deformation extends to $\mathfrak{psu}(2|2) \ltimes \mathbb{R}^3$

• $\mathfrak{psu}_q(2|2)$ invariant *R*-matrix

Beisert and Koroteev '08 Beisert, Galleas and Matsumoto '11



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- $\blacktriangleright S = S_0 R \otimes R$

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- ▶ $\mathfrak{psu}_q(2|2)$ invariant *R*-matrix
- $\triangleright S = S_0 R \otimes R$
- S_0 can be found such that S satisfies crossing
- $\mathfrak{psu}_q(2|2)^2$ invariant S-matrix

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• *S*-matrix is physically *pseudo*-unitary ($S^{\dagger} = B S^{-1} B^{-1}$, *B* Herm.)



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- Kinematics of the model? How are excitations described?
 - ► Short representations labeled by central charges U and V (= q^C) satisfying shortening condition



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 - Parametrized by deformed x^{\pm} variables



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- Kinematics of the model? How are excitations described?
 - ► Short representations labeled by central charges U and V (= q^C) satisfying shortening condition
 - Parametrized by deformed x^{\pm} variables
 - ▶ Natural *definition* of *E* and *p* in terms of *U* and *V*



Parametrizing the fundamental representation

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Conclusion

• Central charges in terms of x^{\pm}

$$U^{2} = \frac{1}{q} \frac{x^{+} + \xi}{x^{-} + \xi}, \quad V^{2} = q \frac{x^{+}}{x^{-}} \frac{x^{-} + \xi}{x^{+} + \xi}$$



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• Then the shortening condition is (equivalent to)

$$\frac{1}{q}\left(x^{+}+\frac{1}{x^{+}}\right)-q\left(x^{-}+\frac{1}{x^{-}}\right)=\left(q-\frac{1}{q}\right)\left(\xi+\frac{1}{\xi}\right)$$

with

$$\xi = -\frac{i}{2} \frac{g(q-q^{-1})}{\sqrt{1 - \frac{g^2}{4}(q-q^{-1})^2}}$$



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Conclusion

• To connect smoothly with string theory (q = 1) we define:

$$V^2 \equiv q^H \,, \ \ U^2 \equiv e^{ip}$$



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• To connect smoothly with string theory (q = 1) we define:

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• Then shortening = deformed string dispersion



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- Then shortening = deformed string dispersion
- $H \to i\tilde{p}, p \to i\tilde{H}$: mirror dispersion ($q = e^{i\pi/k}$)



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Finite size AdS/CFT

The *q*-deformed model

q-deformed spin chain TBA

The q-deformed AdS₅ \times S⁵ mirror TBA

Conclusion

• To connect smoothly with string theory (q = 1) we define:

$$V^2 \equiv q^H, \ U^2 \equiv e^{ip}$$

• Then shortening = deformed string dispersion

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$$H \to i\tilde{p}, p \to i\tilde{H}$$
: mirror dispersion $(q = e^{i\pi/k})$

$$\tilde{H} = 2 \operatorname{arcsinh}\left(\frac{1}{g} \frac{\sin \frac{\pi}{2k}}{\sin \frac{\pi}{k}} \sqrt{1 + \left(1 + g^2 \sin^2 \frac{\pi}{k}\right) \frac{\sinh^2 \frac{\pi}{2k}\tilde{p}}{\sin^2 \frac{\pi}{2k}}}\right)$$



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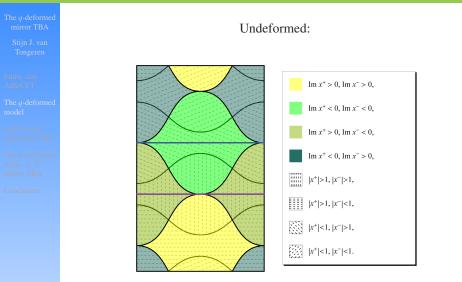
• As for q = 1, this can be uniformized on a torus ("torus = space of short reps")



The dispersion relation on the torus

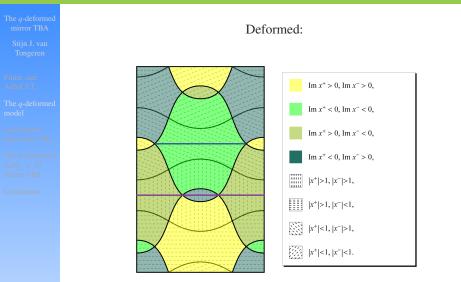


The dispersion relation on the torus





The dispersion relation on the torus





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Conclusion

• Recall $S = S_0 R \otimes R$



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with two sets of

$$1 = \prod_{i=1}^{K^{\mathrm{I}}} \sqrt{q} \frac{y_m - x_i^-}{y_m - x_i^+} \sqrt{\frac{x_i^+}{x_i^-}} \prod_{i=1}^{K^{\mathrm{III}}} \frac{\sinh \frac{\pi g}{2k} \left(v_m - w_i - \frac{i}{g}\right)}{\sinh \frac{\pi g}{2k} \left(w_n - w_i + \frac{i}{g}\right)},$$

$$-1 = \prod_{i=1}^{K^{\mathrm{II}}} \frac{\sinh \frac{\pi g}{2k} \left(w_n - v_i + \frac{i}{g}\right)}{\sinh \frac{\pi g}{2k} \left(w_n - v_i - \frac{i}{g}\right)} \prod_{j=1}^{K^{\mathrm{III}}} \frac{\sinh \frac{\pi g}{2k} \left(w_n - w_j - \frac{2i}{g}\right)}{\sinh \frac{\pi g}{2k} \left(w_n - w_j - \frac{2i}{g}\right)},$$



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• Thermodynamic limit of mABA: string hypothesis



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- Thermodynamic limit of mABA: string hypothesis
 - Physical bound states of the mirror theory (S_0)



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- Thermodynamic limit of mABA: string hypothesis
 - Physical bound states of the mirror theory (S_0)
 - ► String complexes of the auxiliary problem (*R*)





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Conclusion

• Infinite volume mirror theory: bound states?

$$1 = e^{i\bar{p}_l R} \prod_{i \neq l}^{K^{\rm I}} \frac{1}{\sigma^2} \frac{x_l^+ - x_i^-}{x_l^- - x_i^+} \frac{1 - \frac{1}{x_l^- x_i^+}}{1 - \frac{1}{x_r^+ x_i^-}}$$

• $\text{Im}(\tilde{p}_1) > 0$: bound state condition $x_1^- = x_2^+$, multiple solutions



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- Unique solution: physical mirror region



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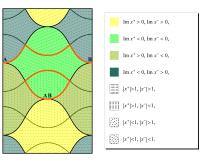
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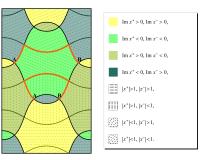
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Conclusion

• Nice parametrization of the physical mirror region?

$$x^{\pm} \to x(u \pm i/g)$$



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• q = 1 mirror region $\longleftrightarrow u$ -plane



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$$q = 1$$
 mirror region $\longleftrightarrow u$ -plane

$$x(u) = \frac{1}{2}(u - i\sqrt{4 - u^2})$$



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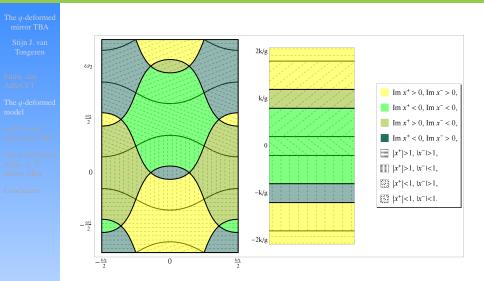
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$$y = e^{i\pi/k} \text{ mirror region} \leftarrow u\text{-plane}$$
$$x(u) = \frac{e^{\frac{\pi gu}{2k}} \left(\sinh\frac{\pi gu}{2k} - i\sqrt{g^2 \sin^2\frac{\pi}{k} - \sinh^2\frac{g\pi u}{2k}}\right) - g^2 \sin^2\frac{\pi}{k}}{g\sin\frac{\pi}{k}\sqrt{1 + g^2 \sin^2\frac{\pi}{k}}}$$







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• Bigger bound states?



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• Bigger bound states?
$$x_1^- = x_2^+, x_2^- = x_3^+, \dots, x_{Q-1}^- = x_Q^+$$



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- Bigger bound states? $x_1^- = x_2^+, x_2^- = x_3^+, \dots, x_{Q-1}^- = x_Q^+$
 - On the *u*-plane we get standard Bethe strings

$$u_j = u + \frac{i}{g}(Q + 1 - 2j), \quad j = 1, \dots, Q$$



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- The deformed theory has a finite spectrum of physical excitations



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- The deformed theory has a finite spectrum of physical excitations
- What about the auxiliary particles?







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Conclusion

• We would like to understand the spectrum associated to *R*





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- We would like to understand the spectrum associated to *R*
- R for $\mathfrak{psu}(2|2) \to \operatorname{Hubbard}$





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- We would like to understand the spectrum associated to *R*
- *R* for $\mathfrak{psu}(2|2) \to \operatorname{Hubbard}$
- *R* for $\mathfrak{psu}_q(2|2) \to q$ -Hubbard?





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- We would like to understand the spectrum associated to R
- *R* for $\mathfrak{psu}(2|2) \to \operatorname{Hubbard}$
- R for $\mathfrak{psu}_q(2|2) \to q$ -Hubbard?
- "Similar" to the q-deformation of the XXX spin chain





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Conclusion

• String hypothesis: Bethe strings of arbitrary length M



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The *q*-deformed AdS₅ \times S⁵ mirror TBA

- String hypothesis: Bethe strings of arbitrary length M
- Y-function for each M-string



The *q*-deformed mirror TBA Stijn J. van Tongeren

Finite size AdS/CFT

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$$\left(s(u) = \frac{1}{4\cosh \pi u/2}\right)$$



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• q-def XXX spin chain is XXZ ($\Delta = \cos \pi/k$)



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 - $M = 1, \dots, k 1$, with $u \in \mathbb{R}$ ("positive parity")
 - M = 1 with Im(u) = ik ("negative parity")



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- Results in special relation: $\tilde{Y}_1 = (Y_{k-1})^{-1}$



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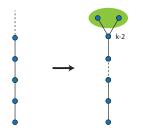
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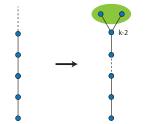
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 $\log Y_{M} = \log (1 + Y_{M+1}) (1 + Y_{M-1}) \star s$ $\log Y_{k-2} = \log (1 + Y_{k-3}) (1 + Y_{k-1})^{2} \star s$ $\log Y_{k-1} = \log (1 + Y_{k-2}) \star s$



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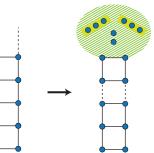
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Saleur and Wehefritz-Kaufmann '00



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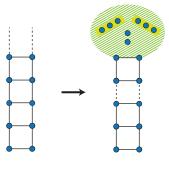
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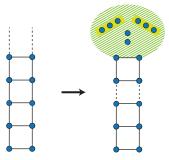
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Saleur and Wehefritz-Kaufmann '00

- $\mathfrak{su}_q(N)$: ???
- $\mathfrak{su}_q(2|2)$: can be nice, elegant, 'simple', real



Quantum deformed Hubbard TBA

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- Come from a pseudo-unitary *R*-matrix $(R^{\dagger} = AR^{-1}A^{-1})$



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 - Self-conjugate spectrum



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- *Mirror* $\mathfrak{psu}_q(2|2)$ string complexes and TBA are 'real'



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Conclusion

• Mirror $\mathfrak{psu}(2|2)$: Hubbard model



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Conclusion

• Mirror $\mathfrak{psu}(2|2)$: Hubbard model (y(v) & w roots)



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The *q*-deformed model

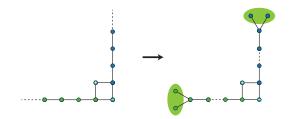
q-deformed spin chain TBA

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Undeformed Mirror TBA



Undeformed Mirror TBA

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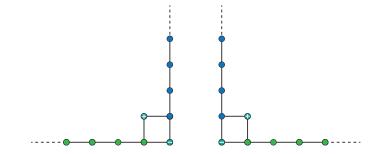
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Conclusion

Two Hubbard subsystems





Undeformed Mirror TBA

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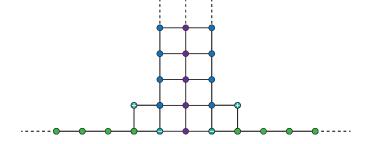
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Conclusion

Two Hubbard subsystems coupled via (∞) *Q*-particles





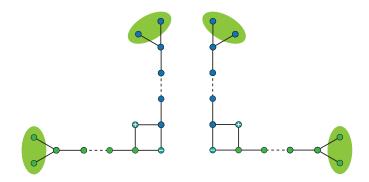
q-deformed Mirror TBA



q-deformed Mirror TBA

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Two q-Hubbard subsystems





q-deformed Mirror TBA

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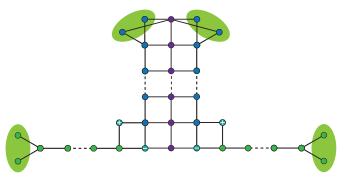
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Two q-Hubbard subsystems coupled via k Q-particles





q-deformed Mirror TBA equations



q-deformed Mirror TBA equations

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$$\begin{split} \log Y_{M|w} &= \log \left(1 + Y_{M+1|w} \right) \left(1 + Y_{M-1|w} \right) \star s - \log \left(1 + Y_{M+1} \right) \star s + \delta_{M,1} \log \left(\frac{1 - Y_{-}}{1 - Y_{+}} \right) \hat{\star} s \\ \log Y_{k-2|w} &= \log \left(1 + Y_{k-3|w} \right) \left(1 + Y_{k-1|w} \right)^{2} \star s - \log \left(1 + Y_{k-1} \right) \star s , \\ \log Y_{k-1|w} &= \log \left(1 + Y_{k-2|w} \right) \star s - \log \left(1 + Y_{k} \right) \star s , \\ \log Y_{M|w} &= \log \left(1 + Y_{M+1|w} \right) \left(1 + Y_{M-1|w} \right) \star s + \delta_{M,1} \log \left(\frac{1 - Y_{-}^{-1}}{1 - Y_{+}^{-1}} \right) \hat{\star} s , \\ \log Y_{k-2|w} &= \log \left(1 + Y_{k-3|w} \right) \left(1 + Y_{k-1|w} \right)^{2} \star s , \\ \log Y_{k-2|w} &= \log \left(1 + Y_{k-3|w} \right) \left(1 + Y_{k-1|w} \right)^{2} \star s , \\ \log Y_{k-1|w} &= \log \left(1 + Y_{k-2|w} \right) \star s , \\ \end{split}$$

$$\log Y_{\pm} = -\log (1+Y_Q) \star K_{\pm}^{Qy} + \log \frac{1+Y_{M|_{W}}^{-1}}{1+Y_{M|_{W}}^{-1}} \star K_M + \log \frac{(1+Y_{k-1|_{VW}})}{(1+Y_{k-1|_{W}})} \star K_{k-1}$$



q-deformed Mirror TBA equations

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$$\begin{split} \log Y_1 &= \log \frac{\left(1 - Y_-^{-1}\right)^2}{1 + Y_2^{-1}} \star s - \check{\Delta} \check{\star} s \,, \\ \log Y_Q &= \log \frac{Y_{Q+1} Y_{Q-1}}{(1 + Y_{Q-1})(1 + Y_{Q+1})} \star s + \log \left(1 + Y_{Q-1}^{-1}|_{vw}\right)^2 \star s \,, \\ \log Y_k &= 2 \log Y_{k-1} \star s - \log (1 + Y_{k-1}) \star s + \log \left(1 + Y_{k-1}^{-1}|_{vw}\right)^4 \star s \,. \end{split}$$



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Conclusion

• The presented TBA equations are in simplified form; closest to Y-system



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- The presented TBA equations are in simplified form; closest to Y-system
- They are derived from so-called canonical equations by applying

$$(K+1)_{MN}^{-1} = \delta_{M,N} - (\delta_{M,N+1} + \delta_{M,N-1})s$$

relying on identities satisfied by kernels for N and $N\pm 1$ bound states



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- For XXZ type equations this still works; would-be length *k* bound states scatter trivially (add zero)

$$Y_{k-1|w}^+ Y_{k-1|w}^- = 1 + Y_{k-2|w}$$



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• For our momentum carrying particles this is *not* the case



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Conclusion

• Still, we derived

$$\log Y_{k} = 2 \log Y_{k-1} \star s - \log(1 + Y_{k-1}) \star s + \log \prod_{\alpha=1,2} \left(1 + \frac{1}{Y_{k-1|w}^{(\alpha)}} \right)^{2} \star s$$



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Finite size AdS/CFT

The *q*-deform model

q-deformed spin chain TBA

The *q*-deformed AdS₅ \times S⁵ mirror TBA

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- Idea: if we had a length k + 1 bound state we would be ok at k
- Nice relation between k + 1 and k 1?

$$S_{k+1}(u) = S_{k-1}(u) \underbrace{S_1(u+ik/g)S_1(u-ik/g)}_{S_1(u-ik/g)}$$



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- For auxiliary kernels the remainder are some known kernels
- For S_0 , precisely with $q = e^{i\pi/k}$ we get crossing!
- Total remainder is then just the equation for $Y_{k-1|vw}$; done



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Conclusion

Reversing the logic



Crossing and the finite Y-system III

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Conclusion

Reversing the logic

• Assuming the bound state S_0 satisfies discrete Laplace

$$\frac{S_{MN}^{+}S_{MN}^{-}}{S_{MN+1}S_{MN-1}} = 1$$



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we can 'derive' the crossing equation!





The <i>q</i> -deformed mirror TBA
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Conclusion

• TBA in finite size AdS/CFT





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- TBA in finite size AdS/CFT
- q-deformed mirror model: spectrum bounded





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 - XXX to XXZ: interesting TBA structure





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- q-deformed mirror TBA and Y-system
 - Closure relies *essentially* on crossing





Finite size AdS/CFT

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Conclusion

• Excited states via asymptotic solution (coming soon)





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The *q*-deformed AdS₅ \times S⁵ mirror TBA

- Excited states via asymptotic solution (coming soon)
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Finite size AdS/CFT

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- Excited states via asymptotic solution (coming soon)
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- Further insight into the $AdS_5 \times S^5$ mirror model



Outlook

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- Deformation with q real
- (TBA for) q-Hubbard proper (Alcaraz-Bariev)

