# The quantum deformed mirror TBA 

Stijn J. van Tongeren

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Work done in collaboration with G. Arutyunov and M. de Leeuw, [1208.3478]

## Introduction



Arutyunov, Frolov '09
Gromov, Kazakov, Vieira '09
Bombardelli, Fioravanti, Tateo '09
Gromov, Kazakov, Kozak, Vieira ’09

## Introduction

The $q$-deformed mirror TBA

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## Motivation

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The $q$-deformed mirror TBA

- Interesting model


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Alcaraz and Bariev '99

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- $q$-deformed theory interpolates


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- $q=1, g$ arbitrary: $\mathrm{AdS}_{5} \times \mathrm{S}^{5}$ string theory


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- Physical 'regularization' of the problem $\left(q=e^{i \pi / k}\right)$
- Wider perspective
- Possible (partial) applications to particular deformed backgrounds


## Outline

Finite size AdS/CFT

The $q$-deformed model and its bound states
$q$-deformed spin chain TBA

The $q$-deformed $\mathrm{AdS}_{5} \times \mathrm{S}^{5}$ mirror TBA

Concluding remarks

## Finite size integrability in $\mathrm{AdS}_{5} / \mathrm{CFT}_{4}$

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## Finite size integrability in $\mathrm{AdS}_{5} / \mathrm{CFT}_{4}$

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Finite size
AdS/CFT
Conclusion

- TBA describes finite size string spectrum via a mirror model

Zamolodchikov '89
Ambjorn, Janik and Kristjansen '05
Arutyunov and Frolov '07

## Finite size integrability in $\mathrm{AdS}_{5} / \mathrm{CFT}_{4}$

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- Deformed model
- Thermodynamic limit: different string hypothesis

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The \(q\)-deformed mirror TBA
- The \(\mathfrak{s u}(2 \mid 2)\) superalgebra in Chevalley-Serre basis \((3 \times E, F, H)\)
\[
\left[H_{i}, H_{j}\right]=0,\left[H_{i}, E / F_{j}\right]= \pm A_{i j} E / F_{j}, \quad\left[E_{i}, F_{j}\right\}=\delta_{i j} D_{i} H_{i}
\]
with
\[
A=\left(\begin{array}{ccc}
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plus (deformed) Serre relations
- We take \(q=e^{i \pi / k}\) with integer \(k>2\)

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- \(q\)-deformation extends to \(\mathfrak{p s u}(2 \mid 2) \ltimes \mathbb{R}^{3}\)
- \(\mathfrak{p s u}_{q}(2 \mid 2)\) invariant \(R\)-matrix

Beisert and Koroteev '08

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- Kinematics of the model? How are excitations described?
- Short representations labeled by central charges $U$ and $V\left(=q^{C}\right)$ satisfying shortening condition
- Parametrized by deformed $x^{ \pm}$variables
- Natural definition of $E$ and $p$ in terms of $U$ and $V$


## Parametrizing the fundamental representation

- Central charges in terms of $x^{ \pm}$

$$
U^{2}=\frac{1}{q} \frac{x^{+}+\xi}{x^{-}+\xi}, \quad V^{2}=q \frac{x^{+}}{x^{-}} \frac{x^{-}+\xi}{x^{+}+\xi}
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- Then the shortening condition is (equivalent to)

$$
\frac{1}{q}\left(x^{+}+\frac{1}{x^{+}}\right)-q\left(x^{-}+\frac{1}{x^{-}}\right)=\left(q-\frac{1}{q}\right)\left(\xi+\frac{1}{\xi}\right)
$$

with

$$
\xi=-\frac{i}{2} \frac{g\left(q-q^{-1}\right)}{\sqrt{1-\frac{g^{2}}{4}\left(q-q^{-1}\right)^{2}}}
$$

## The dispersion relation

The $q$-deformed mirror TBA

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- To connect smoothly with string theory $(q=1)$ we define:

$$
V^{2} \equiv q^{H}, \quad U^{2} \equiv e^{i p}
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\tilde{H}=2 \operatorname{arcsinh}\left(\frac{1}{g} \frac{\sin \frac{\pi}{2 k}}{\sin \frac{\pi}{k}} \sqrt{1+\left(1+g^{2} \sin ^{2} \frac{\pi}{k}\right) \frac{\sinh ^{2} \frac{\pi}{2 k} \tilde{p}}{\sin ^{2} \frac{\pi}{2 k}}}\right)
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- $\operatorname{ssss} G$ connection: rescale $\tilde{H} \rightarrow \frac{\tilde{H}}{g}$ and $\tilde{p} \rightarrow \frac{k}{\pi} \frac{\tilde{p}}{g}$, limit $g \rightarrow \infty$

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\tilde{H}^{2}-\tilde{p}^{2}=\cos ^{-2} \frac{\pi}{2 k}
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- As for $q=1$, this can be uniformized on a torus
("torus $=$ space of short reps")


## The dispersion relation on the torus

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## The dispersion relation on the torus

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The $q$-deformed model

Undeformed:

$\operatorname{Im} x^{+}>0, \operatorname{Im} x^{-}>0$,
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шш
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## The dispersion relation on the torus

The $q$-deformed mirror TBA

Stijn J. van
Tongeren

The $q$-deformed model

## Deformed:


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## The mirror Bethe equations

The $q$-deformed mirror TBA

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$$

with two sets of

$$
\begin{aligned}
1 & =\prod_{i=1}^{K^{\mathrm{I}}} \sqrt{q} \frac{y_{m}-x_{i}^{-}}{y_{m}-x_{i}^{+}} \sqrt{\frac{x_{i}^{+}}{x_{i}^{-}}} \prod_{i=1}^{K^{\mathrm{III}}} \frac{\sinh \frac{\pi g}{2 k}\left(v_{m}-w_{i}-\frac{i}{g}\right)}{\sinh \frac{\pi g}{2 k}\left(v_{m}-w_{i}+\frac{i}{g}\right)} \\
-1 & =\prod_{i=1}^{K^{\mathrm{II}}} \frac{\sinh \frac{\pi g}{2 k}\left(w_{n}-v_{i}+\frac{i}{g}\right)}{\sinh \frac{\pi g}{2 k}\left(w_{n}-v_{i}-\frac{i}{g}\right)} \prod_{j=1}^{K^{\mathrm{III}}} \frac{\sinh \frac{\pi g}{2 k}\left(w_{n}-w_{j}-\frac{2 i}{g}\right)}{\sinh \frac{\pi g}{2 k}\left(w_{n}-w_{j}+\frac{2 i}{g}\right)}
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## Bound states

The $q$-deformed mirror TBA

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## Bound states

- Infinite volume mirror theory: bound states?


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The $q$-deformed
mirror TBA
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## Bound states and the $u$-plane

The $q$-deformed mirror TBA

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## Bound states and the $u$－plane

The $q$－deformed mirror TBA

Stijn J．van Tongeren

The $q$－deformed model


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The $q$-deformed
mirror TBA
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- What about the auxiliary particles?


## Auxiliary spectrum?

The $q$-deformed mirror TBA

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The q-deformed
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- \(R\) for \(\mathfrak{p s u}_{q}(2 \mid 2) \rightarrow q\)-Hubbard?
- "Similar" to the \(q\)-deformation of the XXX spin chain

\section*{TBA for the XXX spin chain}
The \(q\)-deformed
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Finite siza
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\section*{TBA for the XXZ spin chain II}

The \(q\)-deformed mirror TBA
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- \(\mathfrak{s u}_{q}(2 \mid 2)\) : can be nice, elegant, 'simple', real

\section*{Quantum deformed Hubbard TBA}

\author{
The \(q\)-deformed mirror TBA \\ Stijn J. van \\ Tongeren
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- Our model has \(\mathfrak{p s u}_{q}(2 \mid 2)\) mirror auxiliary Bethe equations

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\section*{TBA for mirror \(\mathfrak{s u}_{q}(2 \mid 2)\)}

The \(q\)-deformed mirror TBA
- Mirror \(\mathfrak{p s u}(2 \mid 2):\) Hubbard model

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\section*{Undeformed Mirror TBA}
The \(q\)-deformed
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Two Hubbard subsystems



\section*{Undeformed Mirror TBA}

The \(q\)-deformed
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Two Hubbard subsystems coupled via ( \(\infty\) ) \(Q\)-particles


\section*{\(q\)-deformed Mirror TBA}
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\section*{The \(q\)-deformed} \(\mathrm{AdS}_{5} \times \mathrm{S}^{5}\)

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Two \(q\)-Hubbard subsystems coupled via \(k Q\)-particles


\section*{\(q\)-deformed Mirror TBA equations}
The \(q\)-deformed
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The \(q\)-deformed \(\mathrm{AdS}_{5} \times \mathrm{S}^{5}\) mirror TBA

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The \(q\)-deformed mirror TBA
\[
\log Y_{M \mid v w}=\log \left(1+Y_{M+1 \mid v w}\right)\left(1+Y_{M-1 \mid \nu w}\right) \star s-\log \left(1+Y_{M+1}\right) \star s+\delta_{M, 1} \log \left(\frac{1-Y_{-}}{1-Y_{+}}\right) \hat{\star s}
\]
\[
\log Y_{k-2 \mid v w}=\log \left(1+Y_{k-3 \mid v w}\right)\left(1+Y_{k-1 \mid v w}\right)^{2} \star s-\log \left(1+Y_{k-1}\right) \star s,
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\[
\begin{aligned}
\log Y_{1} & =\log \frac{\left(1-Y_{-}^{-1}\right)^{2}}{1+Y_{2}^{-1}} \star s-\check{\Delta} \check{\star} s, \\
\log Y_{Q} & =\log \frac{Y_{Q+1} Y_{Q-1}}{\left(1+Y_{Q-1}\right)\left(1+Y_{Q+1}\right)} \star s+\log \left(1+Y_{Q-1 \mid \nu w}^{-1}\right)^{2} \star s,
\end{aligned}
\]
\[
\log Y_{k}=2 \log Y_{k-1} \star s-\log \left(1+Y_{k-1}\right) \star s+\log \left(1+Y_{k-1 \mid v w}^{-1}\right)^{4} \star s .
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\section*{Crossing and the finite Y-system}

The \(q\)-deformed mirror TBA
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The \(q\)-deformed mirror TBA

Stijn J. van Tongeren
- Still, we derived
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\log Y_{k}=2 \log Y_{k-1} \star s-\log \left(1+Y_{k-1}\right) \star s+\log \prod_{\alpha=1,2}\left(1+\frac{1}{Y_{k-1 \mid v w}^{(\alpha)}}\right)^{2} \star s
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- Total remainder is then just the equation for \(Y_{k-1 \mid v w}\); done

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The \(q\)-deformed
mirror TBA
Stijn J. van
Tongeren

Reversing the logic

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- Assuming the bound state \(S_{0}\) satisfies discrete Laplace
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\frac{S_{M N}^{+} S_{M N}^{-}}{S_{M N+1} S_{M N-1}}=1
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- and the existence of a Y-system
we can 'derive’ the crossing equation!

The \(q\)-deformed mirror TBA

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- \(q\)-deformed mirror TBA and Y-system
- Closure relies essentially on crossing

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- Deformation with \(q\) real
- (TBA for) \(q\)-Hubbard proper (Alcaraz-Bariev)

The \(q\)-deformed mirror TBA

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